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PROJECT RAND

RESEARCH MEMORANDUM

A NEW MODEL FOR FALLOUT CALCULATIONS (U) R. R. Rapp BM-2115 13 February 1958 Copy No._____

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distributions.

The Zuni rockets passed very nearly over ground zero and therefore were assumed to represent a vertical plane through the diameter of the cloud. In order to further simplify the problem circular symmetry about ground zero will be assumed. The data were then plotted at the appropriate height and distance from ground zero and smooth contours of the reported rocket readings were drawn (Fig. 5).

From the smoothed contours, the total activity in the cloud was computed by a numerical integration. From the total activity and the smooth contours, the fraction of activity per cubic kilofoot was computed as a function of height and radial distance is shown in Table 2.

The estimate of $A_{2}(H)$ was chosen as the marginal distribution which is obtained from the numerical integration of R A(R,H) dR at several different heights. This distribution is shown in Fig. (6). The estimate of $A_{3}(R)$ was chosen as the integration of A(R,H) dH for several different values of R. Figure (7) shows this distribution.

The distribution of activity with particle size was chosen as the distribution presented in reference (1). Accumulative distribution was plotted; points were read from the cumulative function; and numerical differentiation was used to determine values of $\Lambda_1(r)$. Since this is a badly skewed distribution, it is convenient to use as a variable $\ln(r)$ instead of r. Figure (8) shows the activity as a function of size for this distribution.







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		25				2.771	2•45	1.23	.367	•245
		15				1.23	1.528	•562	•306	•18 5
		5					.856	•538	•245	-061
		Height (idt)	ıΛ	15	52	35	54	55	65	52

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Table 2

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Debris Concentration (Fraction per Cubic Kilofoot x 10^5) from Rocket Results as a Function of Radius and Height





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Zuni rocket data

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function of radius of particle

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IV. DISCUSSION

The initial test of the model, which will be referred to as "run 1," was made with the Castle-Bravo wind data taken from Dean and Ohmstead. The details are presented in the Appendix.

The ground position of a few selected particles, which were read from the results of the meteorological part of the computation, are shown in Fig. (9), which also shows the circular standard errors, derived from the discussion of meteorological errors, with the assumption of an error standard deviation, σ_{ϵ} , equal to 5 knots. It appears, from this figure, that <u>random</u> errors of this magnitude will not produce large errors in the orientation and general configuration of the pattern. It does seem reasonable to assume, however, that this type of error may cause a considerable error in the fraction of fallout at some point near the boundary of the fallout area. It must be emphasized that these wind error estimates do not include any gross error in defining the vector wind field.

Several variations of the activity functions were tried with the basic wind pattern of run 1. The purpose of these calculations was to make some qualitative estimates of the effect of changing the distribution functions. The details of these calculations are given in the Appendix and are referred to as run 1a, 1b, etc. The inferences which may be drawn from an inspection and comparison of the results of the calculations will be taken up in order.

The difference between runs la and lb was only in the $A_1(r)$ function. The function used in lb was a log-normal distribution with the same mean and standard deviation as the distribution shown in Fig. (6). The difference between the patterns is not striking. The lb pattern is not as intense at the hot spot as the la pattern, but this was expected since the lb pattern



Fig.9---Ground position and circular standard errors of selected particles

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had less material concentrated near the mean.

A comparison of runs a and b shows several interesting facts. The largest change is about 30 per cent at the point of maximum deposition in both of the patterns. The log-normal distribution of run lb has fewer particles in the size below 55μ than does the RAND distribution of run la. The distribution used in run lb also shows a higher close-in value than does the run la, indicating too much activity on larger particles.

A comparison of both of these runs with the fragmentary data of the Castle-Bravo chot brings two facts to light. The first point is that both the la and lb distributions show far too large a fraction on the islands of the Bikini Atoll. The second point is that values are low in the northern part of Rongelap.

These errors may be due to an incorrect wind analysis, but they could also be explained by a set of incorrect distribution functions. In any event, the relatively slight change in the patterns between an empirical distribution and a mathematical model seems to be slight. This suggests that an attempt could be made to fit the best possible log-normal, or any other useful function, to the existing data.

Runs lc and ld were made to note the comparison of the old model with the new. Run lc was the mushroom portion, and run ld was the stem; together they reproduce a pattern with essentially the same assumptions as the old RAND model. The pattern for run lc, the mushroom portion, again shows relatively slight differences from runs la and lb. The addition of the stem material from run ld makes little change in the pattern. The greatest contribution from the stem is near ground zero, which is already too high, and near the perimeter of the pattern where the uncertainty is greatest.



A rather drastic change in the distribution of activity with size was made in run 1s. The amount of activity on particles greater than 95μ radius was halved, and this excess was added to the particles between 35μ and 65μ in radius. This run used the same spatial distribution as runs 1a and 1b. The "hotspot" for run 1e was moved far away from ground zero and was reduced in intensity.



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V. CONCLUSIONS

Since the observations on Castle-Bravo were fragmentary, there is no useful purpose to be served by continuing variation of parameters on the set of winds from this shot. The work done has proven the feasibility of the approach and has verified the accuracy of the model within the limits of the available data. It is believed that the next set of data processed can proceed on the assumptions that:

1. small changes in the distribution functions will produce only negligible changes in the pattern;

2. the wind patterns are sufficiently accurate to put limits on the possible distribution functions;

3. the gross effects of the fallout from large-yield shots are not affected by stem material.

It may be inferred that most of the material is lodged on particles in a very narrow range of sizes and that it is concentrated into a very narrow range of heights at the time of stabilization.

The next set of runs will be made on the Zuni shot of Operation Redwing. An attempt will be made to adjust the distribution functions so as to produce calculated values of the fraction-down which are within the range of error measurement at those stations for which adequate measurements are available. This then will be assumed to be the optimum model, and the errors of this model will be assumed to be a minimum for fallout calculations.



APPENDIX

A. METFOROLOGICAL DATA

Run 1 refers to those failout calculations that were based on the winds at the time of the Castle-Bravo event. The winds were read from the Dean and Ohmstead charts at 8 locations. Table Al gives these locations in terms of N-8 and E-W distances from ground zero. It was necessary to insure that there were no three points which could be simultaneously colinear and closest to the falling particle, because this situation would invalidate the interpolation functions. The winds were read at two different times (8.25 hours and -15.75 hours from H hour). For times longer than 8.25 hours the wind was assumed to be invariant with time. Table A2 lists the winds at the eight points at two different times.

Table Al

Point	Distance East (n mi)	Distance North (n	mi)
A	- 14	50	
В	136	50	
С	286	100	
D	436	50	
E	- 14	100	
F	136	-100	
G	286	- 50	
Н	436	100	

Positions of Points, Relative to Ground Zero, Where Wind Information was a Machine Input

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Table A2a

Wind Values for Run 1 at 15.75 Hours before Shot Time

V = Speed (kn) $\alpha = Direction (10's of degrees from N)$ 3 13 2 6 2 2 2 2 2 2 2 80,000 75,000

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		Elevation (ft)	5,000	10,000	15,000	20,000	25,000	30 , 000	35,000	1:0°, 0:00	4:5,000	50 , cco	55,000	ố 0, 000	65,000	70,000	75,000	30,000	$\chi = Speed (km)$ $\alpha = Direction (1)$

Table A2b

Wind Values for Run 1 at 3.25 Hours after Shot Time

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B. GEOMETRIC DATA

Run la

In order to adjust the geometric distribution functions A_2 and A_3 to to the yield of the Castle-Bravo event, it is necessary to spread the debris over a larger region of space. Since the integrals of the A functions must equal 1, the actual activities must be reduced accordingly. According to Kellogg⁽¹⁰⁾ the cloud diameter of a 3.5 MT device is ~20 n mi and a 15 MT device is ~ 29 n mi. Since the Zuni distribution of $A_3(R)$ reduced to nearly zero at ~17 n mi (see Fig. 8), it was scaled up to ~24 n mi for the Castle-Bravo shot. The activity A_3 was therefore reduced. In a similar way the height of the cloud was increased, and the activity fraction A_2 was decreased.

All of the values of A were tabulated to such a degree that linear interpolation caused less than 1 per cent in the difference between the curves and the interpolated values. The tabulated distribution functions are shown in Table A3.

The machine output gives the fraction of device per square kilofoot at a number of pre-chosen points. The results of this calculation for run la are shown in Fig. Al. The isopleths that are drawn on Fig. 9, are subjective estimates of the lines of equal fraction down.

Run 1b

The conditions for run]b are the same as for run la except for the distribution of activity with particle size. The distribution used in run lb is the log-normal distribution with the same mean and standard deviation as the distribution in run la. Table Λ_1 shows the values of $A_1(r)$ as a function of r and also the differences between $A_1(r)$ for run la and $A_1(r)$

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Activity Function for Run la

h (kf)	A ₂ (H)	r(µ)	A _] (1nr)	R(n mi)	$A(R) \times 10^{-4}$
0	0 0	9.5	0.15	0	.102
5	.00024	12.25	0.24	.822	.138
10	•00068	1.5.0	0.33	1.645	.174
15	.00140	20.0	0.41	2.467	.211
20	.00240	24.0	0.66	3.289	•500
25	.0037 6	31.0	0.93	4.112	•316
30	•00544	38.0	1.19	4.934	. 385
35	.00700	49.0	1.37	5.757	. 456
40	.01040	60.0	1.42	6.579	.688
45	•01464	77•5	1.41	7.401	•764
50	.02200	95.0	0.68	8.224	•495
55	•035 60	123.0	0.43	9.046	.372
60	•04904	151.0	0.27	9.868	•303
65	•03624	195.0	0.19	10.691	. 252
70	•00280	239.0	0.20	11.513	.218
75	.01 456	309.0	0.16	12.336	•191
80	. 00976	379.0	0.11	13.158	.169
85	•00632	489.0	0.06	13.980	.151
90	.00384	600.0	0	14.803	.131
95	.00208			15.625	.117
100	.00128			16.447	.102
				17.270	.089
				18.092	.076
				18,914	•065
				19.937	•054
	Ì			20.559	•044
	}			21.382	•033
				22.204	.023
				23.026	•014
	Į			23.849	•005



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for run 1b. Figure A2 shows the fallout computation for run 1b with isopleths subjectively drawn.

Table A4

	A _l (r) fo	or Run 16 e	នេ ឧ]	dunet	tion	ı of	(r)),		
and	the,Difference	between A	for	Run	la	and	^_	for	Run	<u>1</u> b

Index	r	A ^b l(r)	$A_1^a - A_1^b$
1	9•5	•280	13
6	12.25	•389	15
11	15.0	•520	19
16	20.0	•648	24
21	24.0	•767	11
26	31.0	. 865	•07
31	38.0	•918	.21
j ó	49 . 0	.928	․ կկ
41	60.0	.883	•54
46	77.5	.800	.61
51	95.0	•685	•00
56	123.0	•559	13
61	151.0	.427	16
66	195.0	•310	12
71	239.0	•214	01
7 6	309.0	.140	.02
81	379.0	. 086	.02
86	489.0	•051	.01
91	600.0	•058	•03

Run 1c

In order to be able to check the effect of the A(R) and A(H) distributions a model was set up which was similar to the model discussed in reference (1). The exponential decrease with height was assumed for the mushroom and the invariant distribution with R was retained. However 97 per cent of the activity was put into the mushroom, instead of 90 per cent as in the earlier model, and the remaining 3 per cent was used to make up run ld. The input values are shown in Table A5 and the final pattern is shown in Fig. (A3).





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Table A5

Activity Distribution for Run le

r	A ₁ (r)	R	А ₃ (R)	h	A ₂ (H)
9.5	0.1455	0	3.513x10 ⁻⁵	0	0
12.25	0.0308	_ ^{Roo}	3-513×10 ⁻⁵	5	O
15.0	0.3201	1.645	3.513x10 ⁻⁵	10	0
20.0	0.3977	2.467	3.513x10 ⁻⁵	15	0
24.0	0.6402	2.467	3.513x10 ⁻⁵	20	0
31.0	0.9021	2.467	3.513x10 ⁻⁵	25	0
38.0	1.1543	2.467	3.513x10 ⁻⁵	30	Û
49.0	1.3289	2.467	3.513x10 ⁻⁵	55	0
60.0	1.3774	2.467	3.513x10 ⁻⁵	40	Û
77.5	1.3677	2.467	3.513x10 ⁻⁵	45	0
95.0	0.6596	2.467	3.513x10 ⁻⁵	50	0
123.0	0.4171	2.467	3.513×10 ⁻⁵	55	0.00560
151.0	0.2619	2.467	3.513×10 ⁻⁵	60	0.04394
195.0	0.1843	2.467	3.513x10 ⁻⁵	65	0.03472
239.0	0.1940	19.737	3.513x10 ⁻⁵	70	0.02743
309.0	0.1552	20.557	0	75	0.02167
379.0	0.1067	20.557	0	80	0.01712
489.0	0.0582	20.557	0	85	0.01353
600.0	0	20.557	0	90	0.01069
		23.849	0	95	0.0084
				100	0.0067

Run 1d

This run was merely the stem portion of run lc. The total fraction placed in the stem was only 3 per cent, and the addition of this fallout to run lc made only minor changes in the pattern. In order to get the gross effects of the fallout, it appears unnecessary to be concerned with the stem. In fact even this small amount of stem fallout seemed to increase the error in the pattern in those areas wherein it was detectable.

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Run 1e

This computation was designed to make a more radical change in the distribution function in order to produce a greater change in the pattern. Since the previous attempts indicated that there was too much debria too close to ground zero, the activity on the larger particles was drastically reduced. The activity size distribution of run le is shown in Table A6 and the resulting pattern is shown in Fig. (A4). It may be noted that the activity is spread further down-wind than in the other runs, although the maximum has been considerably lowered.

Table Aú

The Function $A_1(r)$ for Run le

Index	r	A ₁ (r)
l	9•5	0.15
6	12.25	0.24
11	15.0	0.33
16	20.0	0.41
21	24.0	0.66
26	31.0	0.93
31	38.0	1.59
36	49.0	1.77
41	60.0	1.56
46	. 77.5	1.31
51	95.0	0.34
56	123.0	0.22
61	151.0	0.13
66	195.0	0.10
71	239.0	0.09
7 6	309.0	0.08
81	519.0	0.05
86	489.0	0.03
91	600.0	0



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